Complete Observables for Canonical General Relativity

Bianca Dittrich

Perimeter Institute for Theoretical Physics, Waterloo, Canada

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- GR: almost no observables known ~> approximation?

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 Partial and Complete Observables: method to address these questions

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- Partial and Complete Observables: method to address these questions

Overview

- 1. One Constraint
- 2. Arbitrary Many Constraints
- 3. General Relativity: infinitely many constraints
- 4. Reducing the Number of Constraints: $\infty \rightarrow$ one
- 5. Outlook

Method: C.Rovelli (1990/2002)

Example: FRW-cosmology with massless scalar field

phase space: $\mathbf{x} = (\mathbf{a}, \mathbf{P}_{\mathbf{a}}, \phi, \mathbf{P}_{\phi})$

constraint:
$$C = \frac{1}{2}(-\frac{P_a^2}{a} + \frac{P_{\phi}^2}{a^3})$$

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Complete Observable:
$$\mathbf{F}_{[f; T]}^{\mathcal{T}}(x) = f(x')$$
 where $x' \sim x$ and $T(x') = \tau$

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 $\mathbf{F}_{[a; \phi]}^{\tau = 0}(x) = a \exp(-\operatorname{sgn}(P_a P_{\phi})(0 - \phi))$

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Dirac Observable for arbitrary τ

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 \Rightarrow Idea: Take as many clocks T_i as there are constraints C_i

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Complete Observable: $\mathbf{F}_{[f; T_i]}^{\tau_i}(x) = f(x')$ where $x' \sim x \text{ and } T_i(x') = \tau_i$

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Complete Observable:
$$\mathbf{F}_{[f; T_i]}^{\tau_i}(x) = f(x')$$
 where $x' \sim x$ and $T_i(x') = \tau_i$

 \Rightarrow To compute this introduce a new basis of constraints \tilde{C}_i

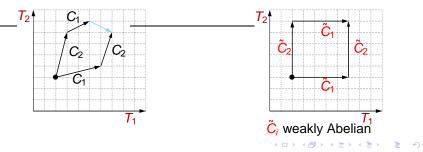
$$\begin{split} \tilde{C}_i &= C_k \, (A^{-1})_{ki} \quad \text{with} \quad A_{lm} = \{ T_l, C_m \} \\ \Rightarrow \quad \{ T_k, \tilde{C}_l \} \simeq \delta_{ki} \quad \Rightarrow \; \{ \tilde{C}_l \} \text{ 'weakly Abelian'} \end{split}$$

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$$\mathbf{F}_{[f; T_i]}^{\tau_i}(x) = \sum_{k_i=0}^{\infty} \frac{1}{|k|!} \tilde{C}_1^{k_1} \cdots \tilde{C}_N^{k_N}[f](x) (\tau_1 - T_1(x))^{k_1} \cdots (\tau_N - T_N(x))^{k_N}$$

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acts on f: $\tilde{C}_i[g] := \{g, \tilde{C}_i\}$

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'distance' of x to $T_i = \tau_i$



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:= f(x') with

x' gauge equivalent to x and

$$T_i(\mathbf{x}') = \tau_i \quad \forall \quad i$$

$$\mathbf{F}_{[f; \mathbf{T}_{i}]}^{\tau_{i}}(x) = \sum_{k_{i}=0}^{\infty} \frac{1}{|k|!} \tilde{C}_{1}^{k_{1}} \cdots \tilde{C}_{N}^{k_{N}}[f](x)(\tau_{1} - \mathbf{T}_{1}(x))^{k_{1}} \cdots (\tau_{N} - \mathbf{T}_{N}(x))^{k_{N}}$$

Some Results

- complete set of Dirac observables
- Poisson brackets between complete observables

($\mathbf{F}_{[\cdot; T_i]}^{\tau_i}$ is a symplecto–morphism onto space of Dirac observables)

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- physical Hamiltonians generate evolution wrt clock variables (Kuchař 1972, Rovelli 1990, Thiemann 2004, B.D 2005)
- partial observables invariant under subset of the constraints: "non-perfect clocks"

Complete Observables for GR (B.D. 2005)

- infinitely many constraints: $C_a(\sigma)$, a = 1, 2, 3; $C_{\perp}(\sigma)$
 - $\sigma \in \Sigma$ (spatial hypersurface)

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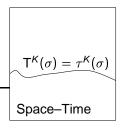
 \rightsquigarrow infinitely many clocks: $T^{\kappa}(\sigma), \ \kappa = 0, \dots 3, \ \sigma \in \Sigma$

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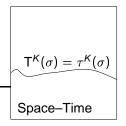
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- $T^{\kappa}(\sigma) = \tau^{\kappa}(\sigma)$ fix embedding of Σ into space-time
- $\mathbf{F}_{[f; \mathbf{T}^{K}(\sigma)]}^{\tau^{K}(\sigma)} = \text{value of } f \text{ on this embedding}$
 - example: $\mathbf{f} = \phi(\sigma^*)$

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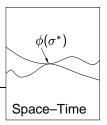
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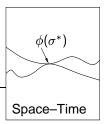
- infinitely many parameters needed to fix embedding
- but: choose f such that it does not depend on all aspects of the embedding
- \rightsquigarrow f invariant under a subset of the constraint: \rightarrow space-time scalar



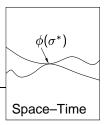
• condition for space-time scalar (K. Kuchař 1977): $\{C[N, N^a], \phi(\sigma^*)\} = 0$ for $N(\sigma^*) = 0 = N^a(\sigma^*)$

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 $\rightsquigarrow \phi(\sigma^*)$ invariant under almost all constraints



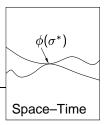
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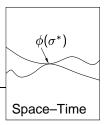
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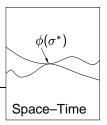


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Need only four constraints for calculation of complete observable

$$\mathbf{F}_{\left[\phi(\sigma^{*}); \mathbf{T}^{K}(\sigma)\right]}^{\tau^{K}(\sigma^{*})} = \sum_{k_{\mathsf{K}}=0}^{\infty} \frac{1}{|\mathbf{k}|!} \tilde{C}_{0}^{k_{0}}[1] \cdots \tilde{C}_{3}^{k_{3}}[1] \left[\phi(\sigma^{*})\right] (\tau_{0} - \mathbf{T}^{0})^{k_{0}} \cdots (\tau_{3} - \mathbf{T}_{3})^{k_{3}}$$

depends on only four parameters $\tau^{K}(\sigma^{*})$

Are there enough Space–Time Scalars?

• scalar matter fields, curvature invariants ~> physical coordinates

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scalar matter fields, curvature invariants → physical coordinates

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•
$$\gamma^{KL} = T^{K}_{\mu}T^{K}_{\nu}\gamma^{\mu\nu} = T^{K}_{,a}T^{L}_{,b}g^{ab} - \{T^{K}, C_{\perp}[1]\}\{T^{L}, C_{\perp}[1]\}$$

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• $\partial_M \gamma^{KL} := \{\gamma^{KL}, \tilde{C}_M[1]\}, \dots$

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• (over–) complete set of Dirac observables : $\mathbf{F}_{\gamma^{KL}}^{\tau^{K}}, \ \mathbf{F}_{\partial_{M}\gamma^{KL}}^{\tau^{K}}, \quad \tau^{1}, \dots, \tau^{3} \in \mathbb{R}^{3}; \tau^{0} \text{ fixed}$

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* physical Hamiltonian: $\{\mathbf{F}_{\gamma KL}^{\tau K}, \mathbf{H}_{M}^{\tau K}(\sigma)\} = \mathbf{F}_{\partial_{M}\gamma KL}^{\tau K} = \frac{\partial}{\partial \tau_{M}} \mathbf{F}_{\gamma KL}^{\tau K}$

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• diffeomorphism invariant constraints $C_{\perp}(\tau^{C}) := \mathbf{D}_{g^{-1/2}C_{\perp}}^{\tau^{C}}$

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(J.Brown, K. Kuchař (1995): another mechanism to get Abelian constraints)

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Need only one constraint for calculation of complete observable!

choose another scalar $T^0(\tau^{\mathbb{C}})$: $\tilde{C}(\tau^{\mathbb{C}}) = [\{T^0(\tau^{\mathbb{C}}), C_{\perp}[1]\}]^{-1}C_{\perp}(\tau^{\mathbb{C}})$

$$\mathbf{F}_{\gamma^{\mathsf{KL}}(\tau^{\mathsf{C}})}^{\tau^{0}} = \sum_{k} \frac{1}{k!} \tilde{C}[1]^{k} [\gamma^{\mathsf{KL}}(\tau^{\mathsf{C}})] (\tau^{0} - T^{0}(\tau^{\mathsf{C}}))^{k}$$

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- * good clock variables? \rightsquigarrow physical Hamiltonian (positive, τ -independent)
- * how to deal with bad clock variables?
- * quantization: different clock variables (with T. Thiemann)

$$\begin{split} C_{\perp}(\tau^{C}) &= \frac{1}{\kappa} g^{-\frac{1}{2}} (\rho_{AB} \rho^{AB} - \frac{1}{2} \rho^{2})(\tau^{C}) - \frac{1}{\kappa} g^{\frac{1}{2}} R(\tau^{C}) + \\ &\quad \frac{1}{2\alpha_{0}} g^{-\frac{1}{2}} (\Pi_{0})^{2} (\tau^{C}) + \frac{1}{2\alpha_{0}} g^{\frac{1}{2}} g^{AB} T^{0}_{,A} T^{0}_{,B} (\tau^{C}) + \\ &\quad \frac{1}{2} g^{-\frac{1}{2}} (\Upsilon) \sum_{A=1}^{3} \frac{1}{\alpha_{A}} (\Pi_{A})^{2} (\tau^{C}) + \frac{1}{2} g^{\frac{1}{2}} (\Upsilon) \sum_{A=1}^{3} \frac{1}{\alpha_{A}} g^{AA} (\tau^{C}) + \\ &\quad g^{\frac{1}{2}} \sum_{K=0}^{3} \frac{1}{\alpha_{K}} V_{(K)} (T^{0} (\tau^{C}), \tau^{C}) \end{split}$$

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