

*From semi-classical LQC
to Friedmann Universe*

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Modified cosmological equations within semi-classical LQC

Modified Friedmann's equation:

Singh, Vandersloot, [gr-qc/0507029]

$$H^2 = \frac{\kappa}{3} \tilde{\rho} \left(1 - \frac{\tilde{\rho}}{\rho_c} \right), \quad \tilde{\rho} = \frac{1}{2} \frac{\psi^2}{D} + V(\phi), \quad H = \frac{d \ln a}{dt}, \quad \kappa = 8\pi G,$$

Modified Klein-Gordon equation:

$$\rho_c = \frac{3}{\kappa(\mu_0 \gamma a)^2}, \quad \mu_0 = \sqrt{3}/4, \quad \gamma = 0.2375\dots,$$

$$\dot{\psi} = -3H\psi \left(1 - \frac{1}{3} \frac{d \ln D}{d \ln a} \right) - D \frac{\partial V}{\partial \phi}, \quad \psi = \dot{\phi} \equiv \frac{d\phi}{dt},$$

$$D(q) = \left(\frac{8}{77} \right)^6 q^{3/2} \left\{ 7 \left[(q+1)^{11/4} - |q-1|^{11/4} \right] - 11q \left[(q+1)^{7/4} - \text{sign}(q-1) |q-1|^{7/4} \right] \right\}^6,$$

$$q = \left(\frac{a}{a_*} \right)^2, \quad a_* = \left(\frac{\kappa \gamma \mu_0 J}{3} \right)^{1/2}.$$

Domain of validity

Vandersloot, [gr-qc/0502082]

These equations are valid if $\mu_0 c \ll 1$,

In the classical limit $c = \gamma \dot{a}$,

So that we have: $\mu_0 \gamma \dot{a} \ll 1$,

This is still valid for large volumes!

$$\mu \gg 2J\mu_0.$$

Reduction to the form of Friedmann's equations

Modified cosmological equations can be written in the Friedmann's form:

$$H^2 = \frac{\kappa}{3}\rho_{eff}, \quad \frac{d\rho_{eff}}{dt} + 3H(\rho_{eff} + p_{eff}) = 0,$$

with the following notation:

$$\rho_{eff} = \tilde{\rho} \left(1 - \frac{\tilde{\rho}}{\rho_c} \right), \quad \left\{ \begin{array}{l} \tilde{\rho} = \frac{1}{2} \frac{\psi^2}{D} + V(\phi), \\ \tilde{p} = \frac{1}{2} \frac{\psi^2}{D} \left(1 - \frac{1}{3} \frac{d \ln D}{d \ln a} \right) - V. \end{array} \right.$$

Lidsey, e.a., PRD70(2004)063521

Some implications of modified Friedmann's equations

1. Constraints on the energy density:

$$H^2 = \frac{\kappa}{3} \tilde{\rho} \left(1 - \frac{\tilde{\rho}}{\rho_c} \right) \longrightarrow \boxed{\tilde{\rho} \leq \rho_c}$$

2. Solution for the energy density:

$$\rho_{\pm} = \frac{\rho_c}{2} \left(1 \pm \sqrt{1 - \frac{12H^2}{\kappa\rho_c}} \right) \longrightarrow \boxed{-\sqrt{\frac{\kappa\rho_c}{12}} \leq H \leq \sqrt{\frac{\kappa\rho_c}{12}}}$$

$$0 \leq \rho_- \leq \frac{\rho_c}{2}$$

$$\frac{\rho_c}{2} \leq \rho_+ \leq \rho_c,$$

3. Hubble parameter equation:

$$\dot{H} = -\kappa \frac{\psi^2}{2} \left(1 - \frac{2\tilde{\rho}}{\rho_c} \right) \left(1 - \frac{1}{6} \frac{d \ln D}{d \ln a} \right) + \frac{\kappa}{6H} \dot{\rho}_c \left(\frac{\tilde{\rho}}{\rho_c} \right)$$

Effective theories within s.c. LQC

$$\left\{ \begin{array}{l}
 H^2 = \frac{\kappa}{3} \tilde{\rho} \left(1 - \frac{\tilde{\rho}}{\rho_c} \right), \\
 \dot{H} = -\kappa \frac{\psi^2}{2} \left(1 - \frac{2\tilde{\rho}}{\rho_c} \right) \left(1 - \frac{1}{6} \frac{d \ln D}{d \ln a} \right) + \frac{\kappa}{6H^2} \rho_c \left(\frac{\tilde{\rho}}{\rho_c} \right) \\
 \dot{\psi} = 3H\psi \left(1 - \frac{1}{3} \frac{d \ln D}{d \ln a} \right) + D \frac{\partial V}{\partial \phi},
 \end{array} \right.$$

$a \ll a_*,$
 $\rho \simeq \rho_c$
(ET-II)

$a \ll a_*,$
 $\rho \ll \rho_c$

(ET-I)

$a \gg a_*,$
 $\rho \simeq \rho_c$

(ET-III)

Classical limit: $a \gg a_*,$
 $\rho \ll \rho_c$

Massless scalar field within ET-I

Vereshchagin, JCAP, 07 (2004) 013

Equations:

$$H^2 + \frac{k}{a^2} = \frac{8\pi}{3M_p^2} \left[\frac{\dot{\psi}^2}{2D} + V \right],$$

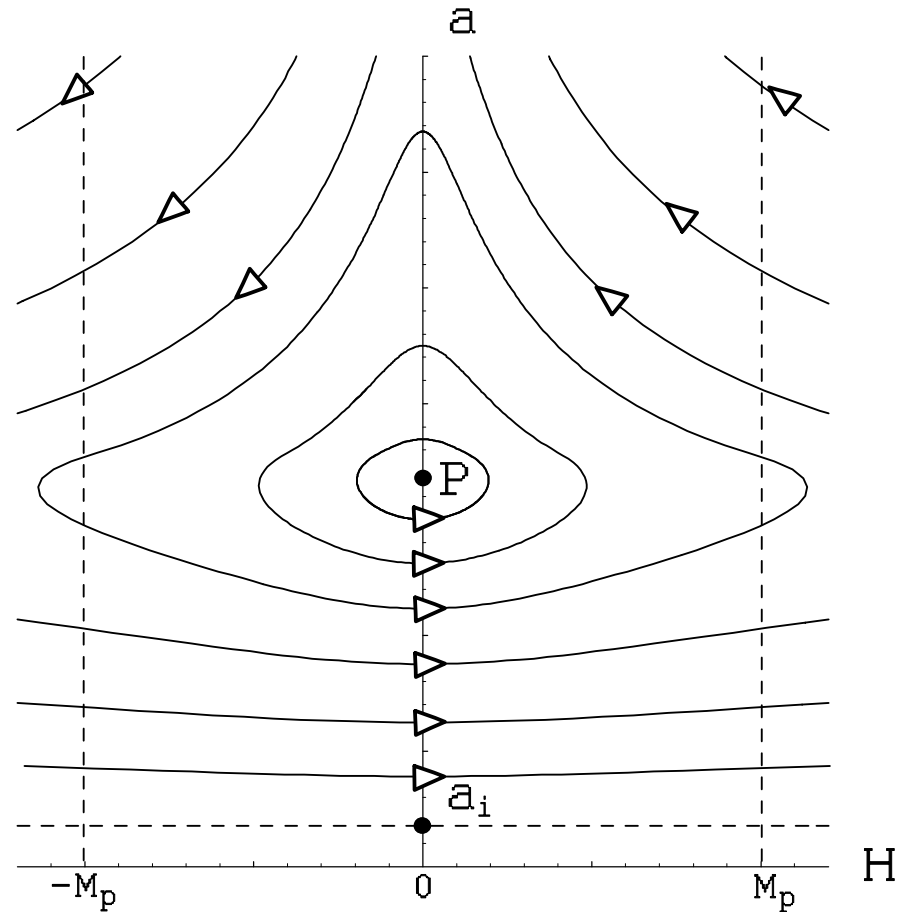
$$\dot{\psi} = -3H\psi + \frac{\dot{D}}{D}\psi - D \frac{\partial V}{\partial \varphi},$$

$$\dot{H} + H^2 = -\frac{8\pi}{3M_p^2} \frac{\dot{\psi}^2}{D} \left[1 - \frac{\dot{D}}{4HD} \right] + \frac{8\pi V}{3M_p^2},$$

For $V=0$ the phase space is reduced to 2D one $\{H, a\}$

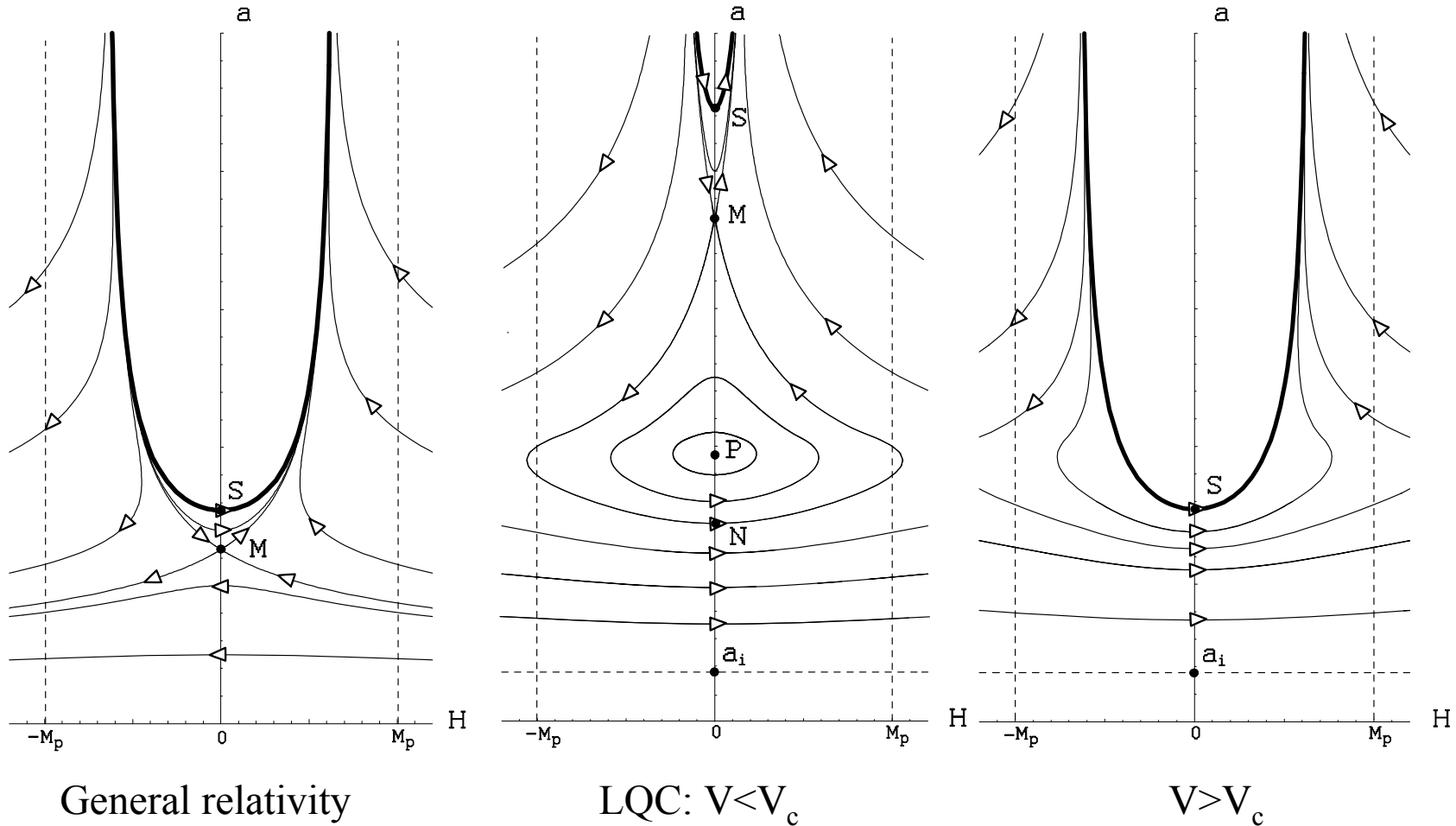
There is only one singular point: center P.

Singularity is absent for all phase trajectories within bounds –
 $M_p < H < M_p$



Massive scalar field within ET-I

Vereshchagin, JCAP, 07 (2004) 013



M – saddle, P – center, S – limiting value of the scale factor.

Singularity avoidance for k=0

$$H^2 = \frac{\kappa}{3} \tilde{\rho} \left(1 - \frac{\tilde{\rho}}{\rho_c} \right), \quad H(0) = 0 \quad \longrightarrow \quad \rho(0) = \rho_c, \quad \text{otherwise} \\ \phi(0) = \dot{\phi}(0) = 0$$

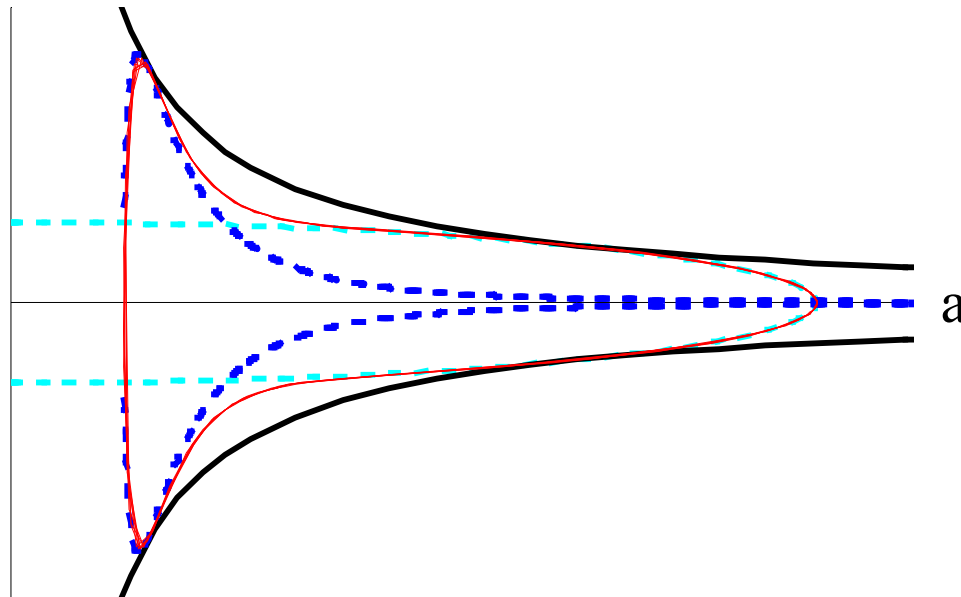
$$\dot{H} = -\kappa \frac{\psi^2}{2} \left(1 - \frac{2\tilde{\rho}}{\rho_c} \right) \left(1 - \frac{1}{6} \frac{a \ln D}{d \ln a} \right) - \frac{\kappa a^2}{\alpha} \tilde{\rho}^2$$

ET-III: $\left\{ \begin{array}{l} H^2 = \frac{\kappa}{3} \left(\frac{1}{2} \psi^2 + V \right) \left[1 - \frac{a^2}{\alpha} \left(\frac{1}{2} \psi^2 + V \right) \right], \\ \dot{H} = -\kappa \frac{\psi^2}{2} \left[1 - \frac{2a^2}{\alpha} \left(\frac{1}{2} \psi^2 + V \right) \right] - \frac{\kappa a^2}{\alpha} \left(\frac{1}{2} \psi^2 + V \right)^2, \\ \dot{\psi} + 3H\psi + D \frac{\partial V}{\partial \phi} = 0. \end{array} \right.$

Cycles within ET-III

- Maximum expansion point is *potential-dominated*
- Bounce is *kinetic-dominated*
- The presence of boundaries is essential

H



Blue – kinetic-dominated solution,

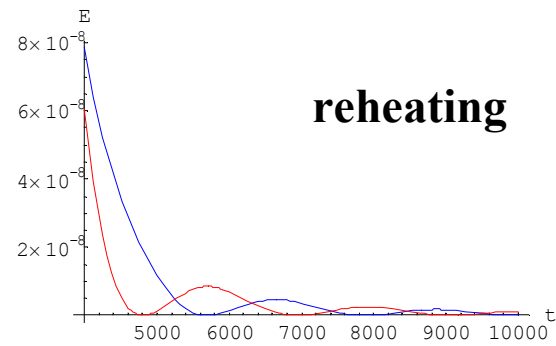
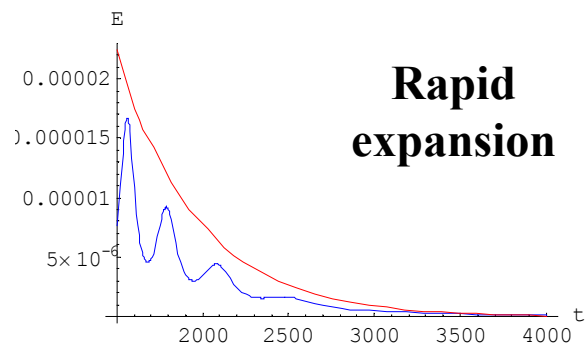
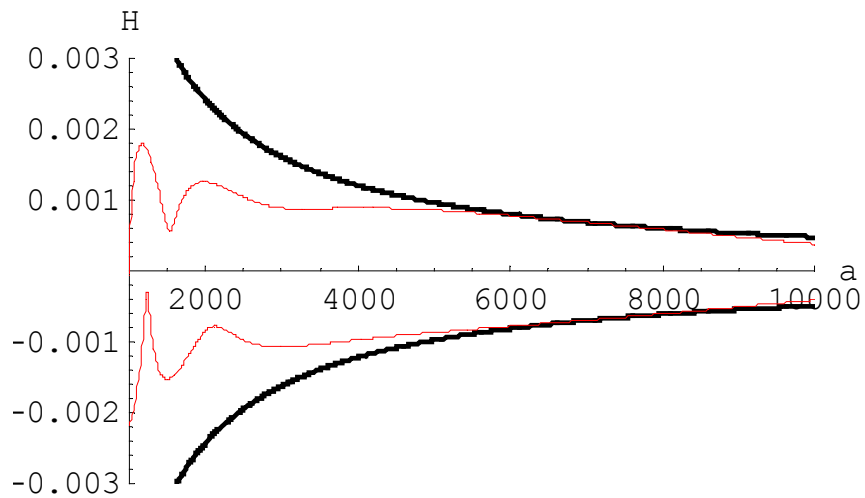
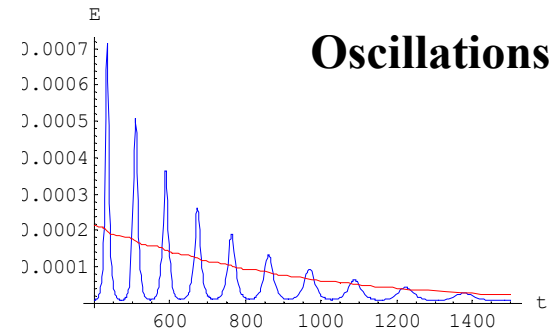
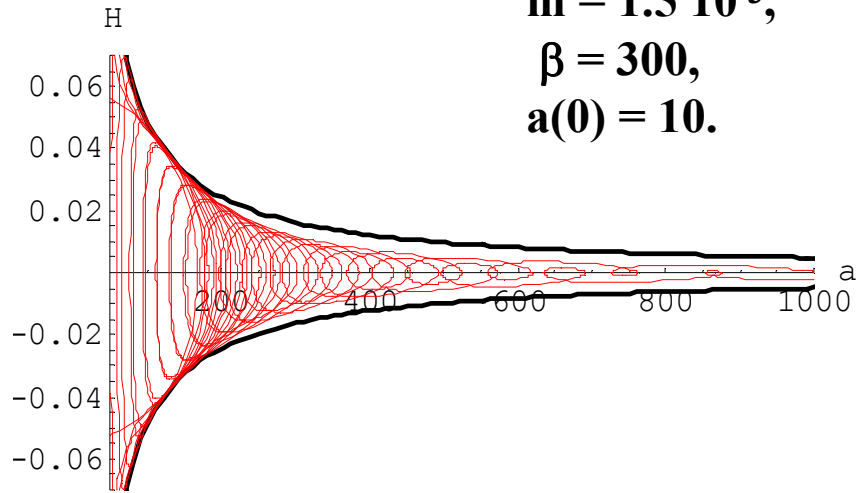
Azure – potential-dominated solution,

Red – numerical solution,

Black - boundaries

Example

$m = 1.5 \cdot 10^{-3}$,
 $\beta = 300$,
 $a(0) = 10$.



Size of the Universe?...

$$\rho_c = \frac{\alpha}{a^2}, \quad \alpha = \frac{3}{\kappa(\mu_0\gamma)^2} \approx 11.28G^{-1},$$

$$Ha \leq \sqrt{\frac{\kappa\alpha}{12}} \approx 5 \quad \longrightarrow \quad a_{today} \leq 5H_0^{-1}$$

... is of the order of $O(H^{-1})$?!

Validity: $Ha \ll (\mu_0\gamma)^{-1} \approx 10$.

Cosmological eras

- **Boundary:** $H \propto a^{-1}$,
- Oscillations: $H \propto a^{-4}$,
- “Potential-dominated”: $H \approx \text{const}$,
- Reheating: $H \propto a^{-3/2}$,
- Radiation-dominated: $H \propto a^{-2}$,
- Matter-dominated: $H \propto a^{-3/2}$,
- Acceleration: $H \approx \text{const}$.

Conclusions and perspectives

- Significance of quantum corrections for the scale factor and/or energy density \rightarrow effective theories;
- ET-III outcomes: absence of singularity, size of the Universe, initial conditions?
- Present day acceleration...close to maximal expansion?
- What about $k=1$ models?
- What is the picture for ET-II?