

Quantum Black Holes I

O. Winkler

Univ. of New Brunswick

CQG 22, L127

CQG 22, L135

PRD 71, 104001

with V. Husain

Hawking radiation

formation

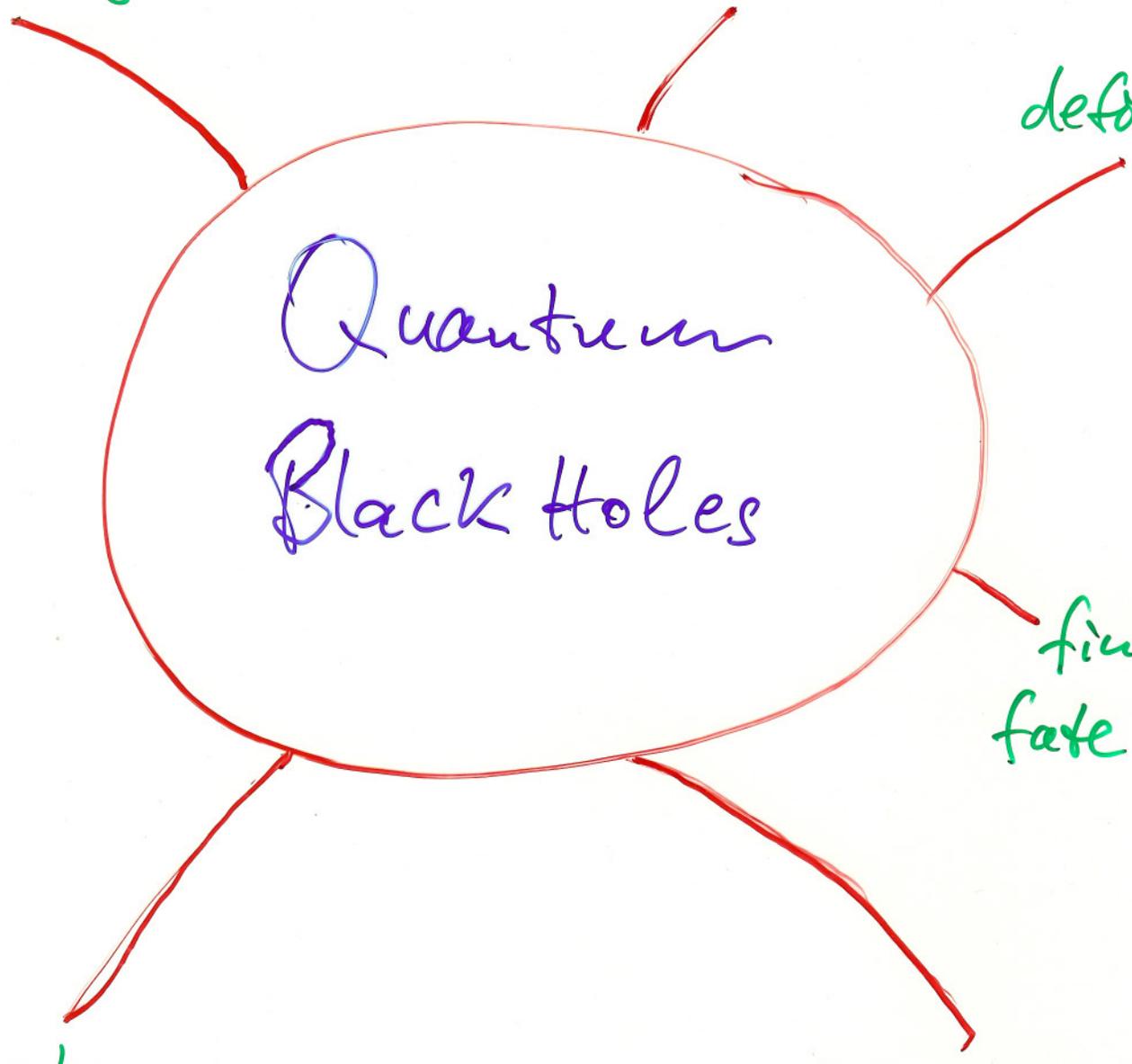
definition

Quantum
Black Holes

final
fate

entropy

Singularity
resolution



Standard model

Spherically symmetric gravity
coupled to matter (scalar field usually)

$$ds_{\Sigma}^2 = \lambda^2(r,t) d\mathbb{R}^2 + R^2(r,t) d\Omega$$

λ, R scalar fields Φ

→ 2D field theory midisuperspace

Classically

System has sectors with BH formation
and sectors w/o bhs (cf. Choptuik)

Task

Quantize the system

Come up with good def. of Quantum BH

Investigate collapse problem

Evaporation?

⋮

70s BCNM
Unruh

fully gauge fixed $\rightarrow \Phi$ only d.o.f.

$$H_{\text{red}} = \int dr \left(\frac{P_\phi^2}{4r^2} + r^2 \phi'^2 \right) \exp\left(\int_\infty^r S_\phi(r') dr' \right)$$

completely untractable

90s Kastrop, Thiemann
Kuchař

fully constrained system
difficult to make contact with
standard semi-cl. calculations like
Hawking radiation



our strategy: partial gauge fixing $\Lambda=1$

R, Φ remaining d.o.f.

Classical Setting

$$q_{ab} = u_a u_b + \frac{R^2}{r^2} (e_{ab} - u_a u_b)$$

$$\tilde{\pi}^{ab} = \frac{P_R}{2} u^{ab} + \frac{r^2 P_R}{4R} (e^{ab} - u^a u^b)$$

(R, P_R) (Φ, P_Φ) phase space var.

fall off conditions chosen by demanding

Painleve-Gullstrand form of metric
in fully gauge fixed form

$$ds^2 = -dt^2 + (dr + \sqrt{\frac{2M}{r}} dt)^2 + r^2 d\Omega^2$$

regular at $r=2M$, covers entire
spatial hypersurface

$$R = r + \mathcal{O}(r^{-1/2-\epsilon})$$

$$P_R = A \frac{r^{-1/2}}{2} + \mathcal{O}(r^{-1-\epsilon})$$

$$\Phi = B r^{-1/2} + \mathcal{O}(r^{-3/2-\epsilon})$$

$$P_\Phi = C r^{1/2} + \mathcal{O}(r^{-\epsilon})$$

Quantization

Methods inspired by LQG

Ex point particle

$$(x, p) \quad H = \frac{p^2}{2m} + V(x)$$

basis states $|e^{ixp}\rangle \equiv |x\rangle$

$$\langle x | y \rangle = \delta_{x,y} \quad \text{Kronecker delta}$$

→ plane waves normalizable
(necessary for singularity resolution)

$$\hat{x}|x\rangle = x|x\rangle$$

\hat{p} does not exist

"non-regular" repr.

only $\hat{u}_a = e^{iap}$

$$\hat{u}_a |x\rangle = |x-a\rangle$$

$$\hat{H}_a = \frac{1}{2m} \frac{z - \hat{u}_a - \hat{u}_a \cdot z}{a^2} + \hat{V}(x)$$

Scalar field = "one particle
at each point $x \in \mathbb{R}_0^+$ "

basis states

$$|e^{i \sum a_k P_R(x_k)}, e^{i \sum b_l P_L(x_l)}, e^{i \sum c_m P_\phi(x_m)}\rangle$$

$$\equiv |a_1 \dots a_{N_1}, b_1 \dots b_{N_2}, c_1 \dots c_{N_3}\rangle$$

$$\langle a_1 \dots a_{N_1}, b_1 \dots b_{N_2}, c_1 \dots c_{N_3} | a'_1 \dots a'_{N_1}, b'_1 \dots b'_{N_2}, c'_1 \dots c'_{N_3} \rangle = \delta_{a_1, a'_1} \dots \delta_{c_{N_3}, c'_{N_3}}$$

- basis states like mode states in Fock space, but in x -space
- represent excitations of quantum field at locations $\{x_k\}$

basic operators

$$R_f = \int_0^\infty dr R(r) f(r)$$



$$\hat{R}_f |a_1 \dots a_N\rangle = \mathcal{Z} \int \sum_k a_k f(x_k) |a_1 \dots a_N\rangle$$

$$e^{i \lambda_j P_R(x_j)} |a_1 \dots a_N\rangle = |a_1 \dots a_j - \lambda_j \dots a_N\rangle$$

Creates excitation λ_j
at location x_j

Singularity Resolution

The "building block", contained in quantities of interest, that diverges at classical singularity is $\frac{1}{R}$

ex. $\tilde{\pi} = \frac{1}{2} \left(\frac{P_A}{R^2} + \frac{P_R}{R} \right)$

H_{red} reduced Hamiltonian

H Hamiltonian constraint

→ strategy: construct $\frac{1}{R}$ operator which has finite eigenvalue at class. singularity

$$\frac{1}{|R_f|} = \left(A e^{-iP_R} \left\{ \sqrt{|R_f|}, e^{iP_R} \right\} \right)^2 \quad \text{Thiemann trick}$$

→ $\widehat{\frac{1}{|R_f|}}$ by "heat operation"

Eigenvalue of \hat{L}^2 at class. singularity

$\frac{z}{l_p^2}$ large but finite!

→ Kinematical resolution of class.
singularity

also dynamics remains finite
at quantum level

What is a Quantum Black Hole?

BHs are intimately connected to the notion of horizon.

Many types of horizon:

Event, apparent, trapping
isolated,

Which one to take?

Want s.th. physically intuitive connected to trapping of light
practical, i.e. useful for "real life"
e.g. numerical relativity

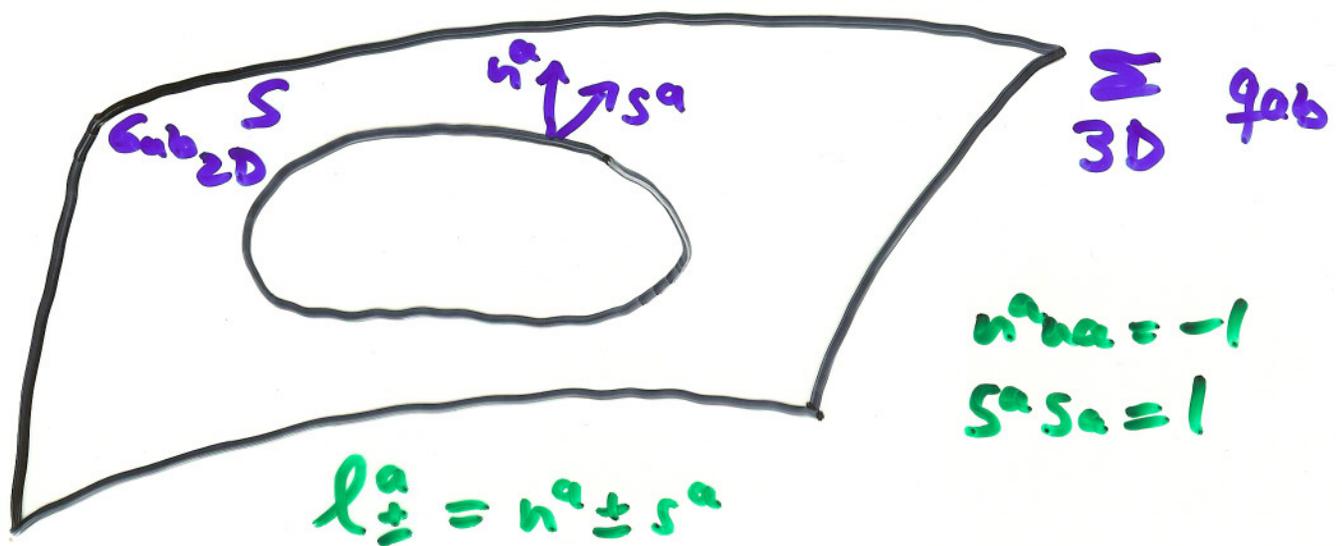
genuinely quantum, i.e. no fixing of classical boundaries

→ trapping horizons

Definition of Quantum Black Holes

Idea

Define BHs via trapped surfaces
(Hayward)



$$\theta_{\pm} = g^{ab} D_a l_{\pm}^b$$

expansion of ^{out}going light rays

S is outer trapped if $\theta_+ \leq 0$ $\theta_- < 0$

marginally outer
trapped

(MOTS)

$$\theta_+ = 0 \quad \theta_- < 0$$

In this framework:

$$\theta_{\pm} = -\frac{1}{2} (4RR' \pm P_{\Lambda})$$

Quantization:

\hat{R}

\hat{P}_{Λ} (see talk by Vigar)

\hat{R}' finite difference of R operators

$\Rightarrow \hat{\theta}_{\pm}$ horizon operator

tool to probe whether state $|\psi\rangle$ describes QBH or not:

if there is \bar{r} such that

$$\langle \psi | \hat{\theta}_{+}(\bar{r}) | \psi \rangle = 0$$

$$\text{and } \langle \psi | \hat{\theta}_{-}(\bar{r}) | \psi \rangle < 0$$

then $|\psi\rangle$ describes quantum spacetime with a qbh with horizon at \bar{r}

$$\text{Horizon size: } R_{H} = \langle \psi | \hat{R} | \psi \rangle$$

Conclusion

- Have kinematical quantum theory
- Singularity resolution
- genuinely quantum definition of a black hole

Next step:

Implement dynamics

↔ next talk